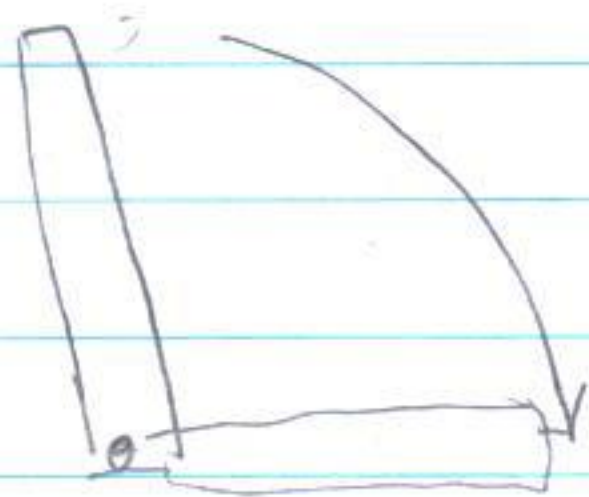


Review Prob

A rod, length L , mass M P61, Ch 10



a) Find its angular speed

$$W_{\text{ext}} = \Delta KE + \Delta PE$$

$$KE_i + PE_i = KE_f + PE_f$$

$$Mg \frac{L}{2} = \frac{1}{2} I \omega_f^2$$

$$I = \frac{1}{3} ML^2$$

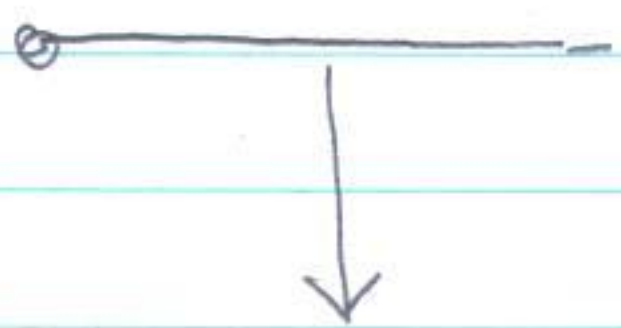
$$Mg \frac{L}{2} = \frac{1}{2} \frac{1}{3} ML^2 \omega_f^2$$

$$\sqrt{\frac{3g}{L}} = \omega_f$$

b) Find the angular acceleration when horizontal

c) Find the x and y components of the \vec{a}_{cm}

d) Determine the force on the pin



$$b) \quad \vec{\tau} = I \vec{\alpha}$$

$$\frac{L}{2} mg \hat{-k} = \frac{1}{3} mL^2 \vec{\alpha}$$

$$\frac{3}{2} \frac{g}{L} \hat{-k} = \vec{\alpha}$$

$$|\vec{\alpha}| = \alpha = \frac{3}{2} \frac{g}{L} \text{ clockwise}$$

$$c) \quad a_{cm}^x = a_c = \omega^2 r_{cm} = \frac{3g}{2L} \cdot \frac{L}{2} = \frac{3}{4}g \text{ in}$$

$$a_{cm}^y = a_t = \alpha r_{cm} = \frac{3}{2} \frac{g}{L} \cdot \frac{L}{2} = \frac{3}{4}g \text{ down}$$

$$d) \quad \sum \vec{F} = m \vec{a}$$

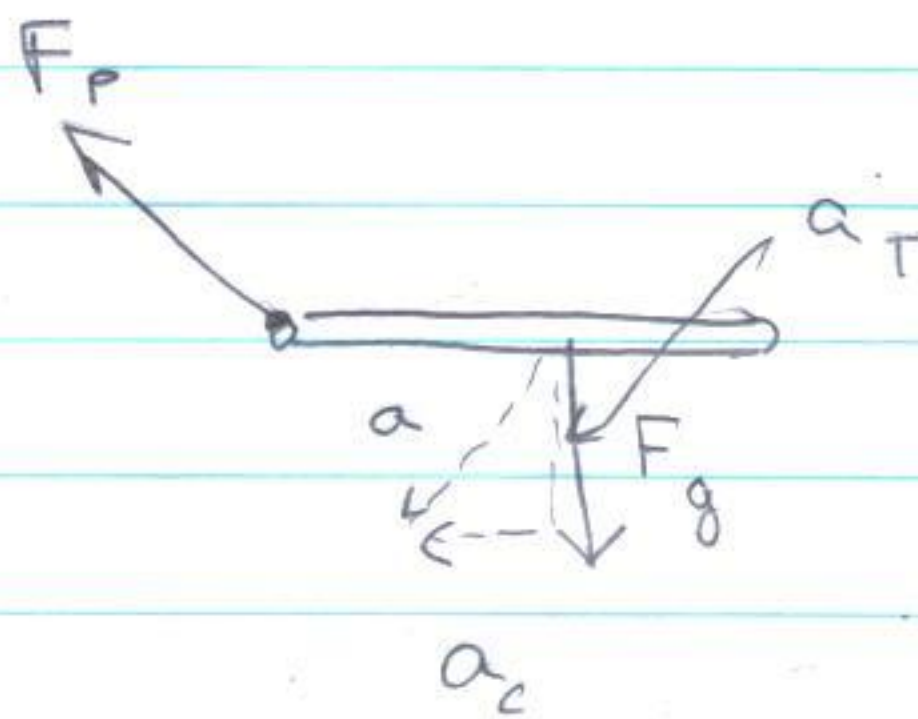
$$\sum F^x = m a^x$$

$$F_p^x = -m a_c$$

$$F_p^x = -m \frac{3}{4}g \leftarrow \text{in}$$

$$\sum F^y = m a^y$$

$$-mg + F_p^y = m \left(-\frac{3}{4}g \right) \Rightarrow F_p^y = \frac{1}{4}mg$$



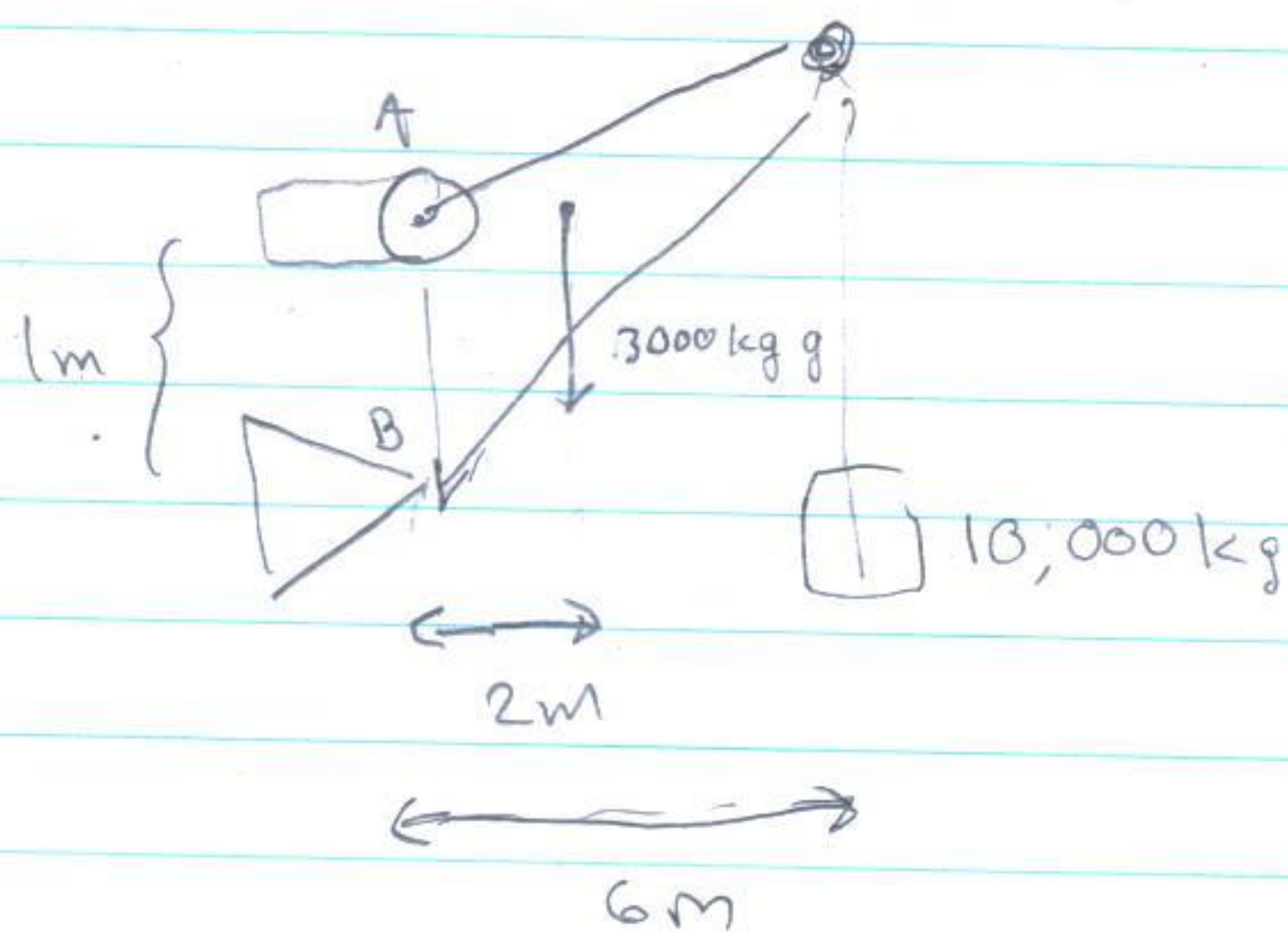
What we have labelled F_p
is the force on board by pin

want $\vec{F}_{\text{on pin}} = -\vec{F}_{\text{board}}$

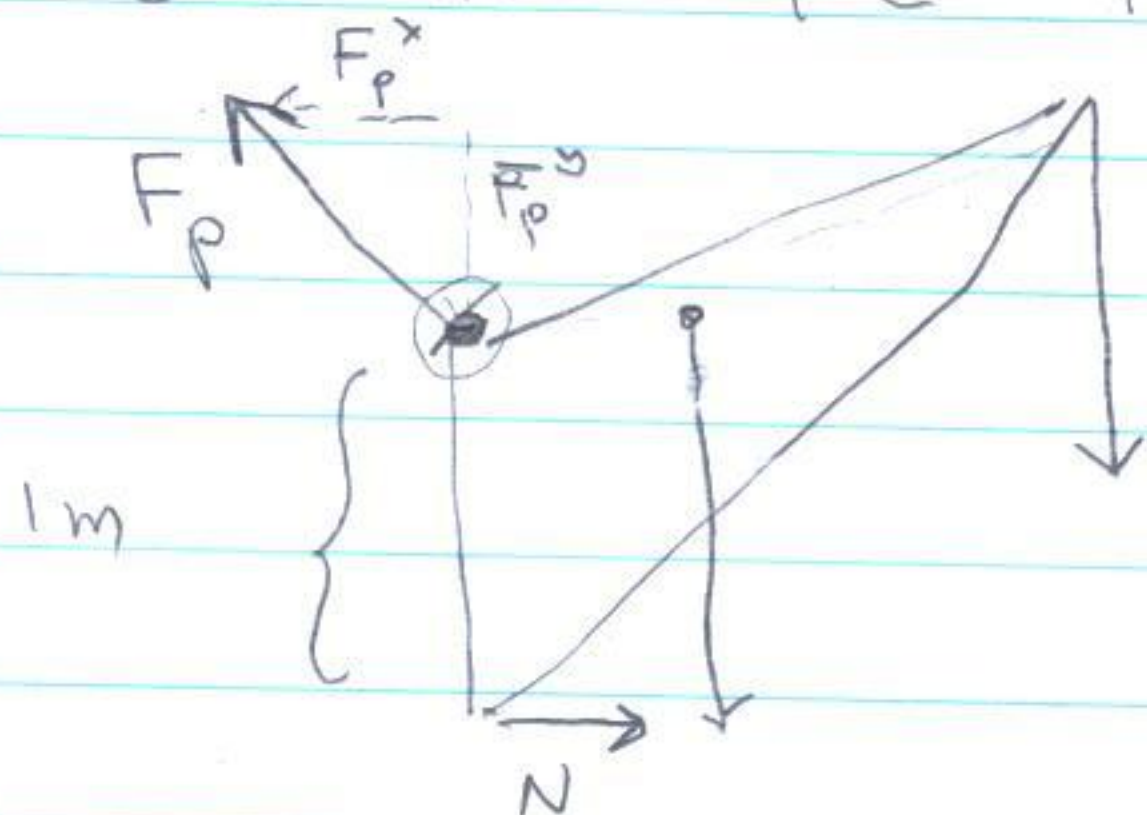
$$\vec{F}_{\text{on pin}} = m \frac{3}{2} g \hat{i} - \frac{1}{4} mg \hat{j}$$

$$|F_p| = mg \left(\left(\frac{3}{2}\right)^2 + \left(\frac{1}{4}\right)^2 \right)^{\frac{1}{2}} \quad \text{etc}$$

Review Prob P47 Ch 12 book



Determine the reaction forces at A & B



$$\sum F^x = 0 \quad \sum F^y = 0$$

$$-F_p^x + N = 0$$

$$-3M_1g - M_2g + F_p^y = 0$$

$$\sum \tau = I \alpha$$

$$+ N(1m) - M_1g(2m) - M_2g(6m) = I \alpha$$

$$|\tau| = |R F_{\perp}| = |R F \sin \theta| = |R_{\perp} F|$$

Spring + Oscillatory Motion



$$F = -kx = ma$$

$$-\frac{k}{m}x = \frac{d^2x}{dt^2} \Rightarrow \frac{d^2x}{dt^2} = -\omega_0^2 x$$

Solution

$$x = A \cos(\omega_0 t + \phi)$$

$$\omega_0 = \sqrt{\frac{k}{m}} = \text{"angular frequency"}$$

implicitly assume
radian measure

= how many
radians
turned per
sec

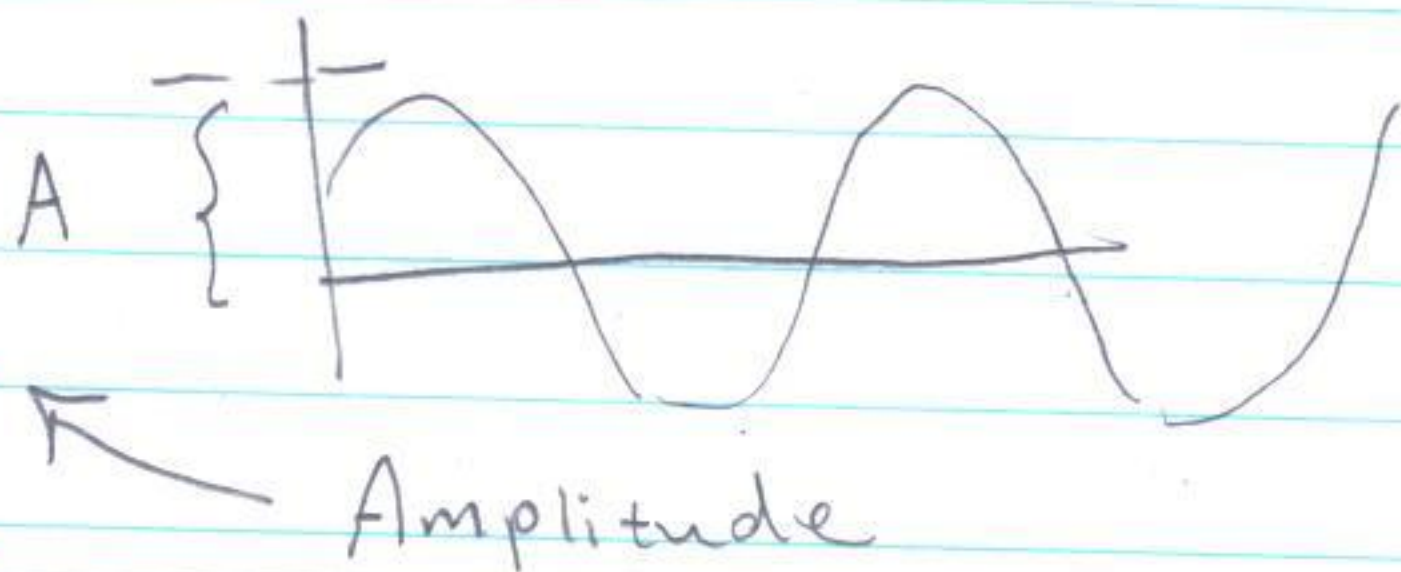
$$v = \frac{dx}{dt} = -A\omega_0 \sin(\omega_0 t + \phi)$$

$$a = \frac{dv}{dt} = -\omega_0^2 [A \cos(\omega_0 t + \phi)]$$

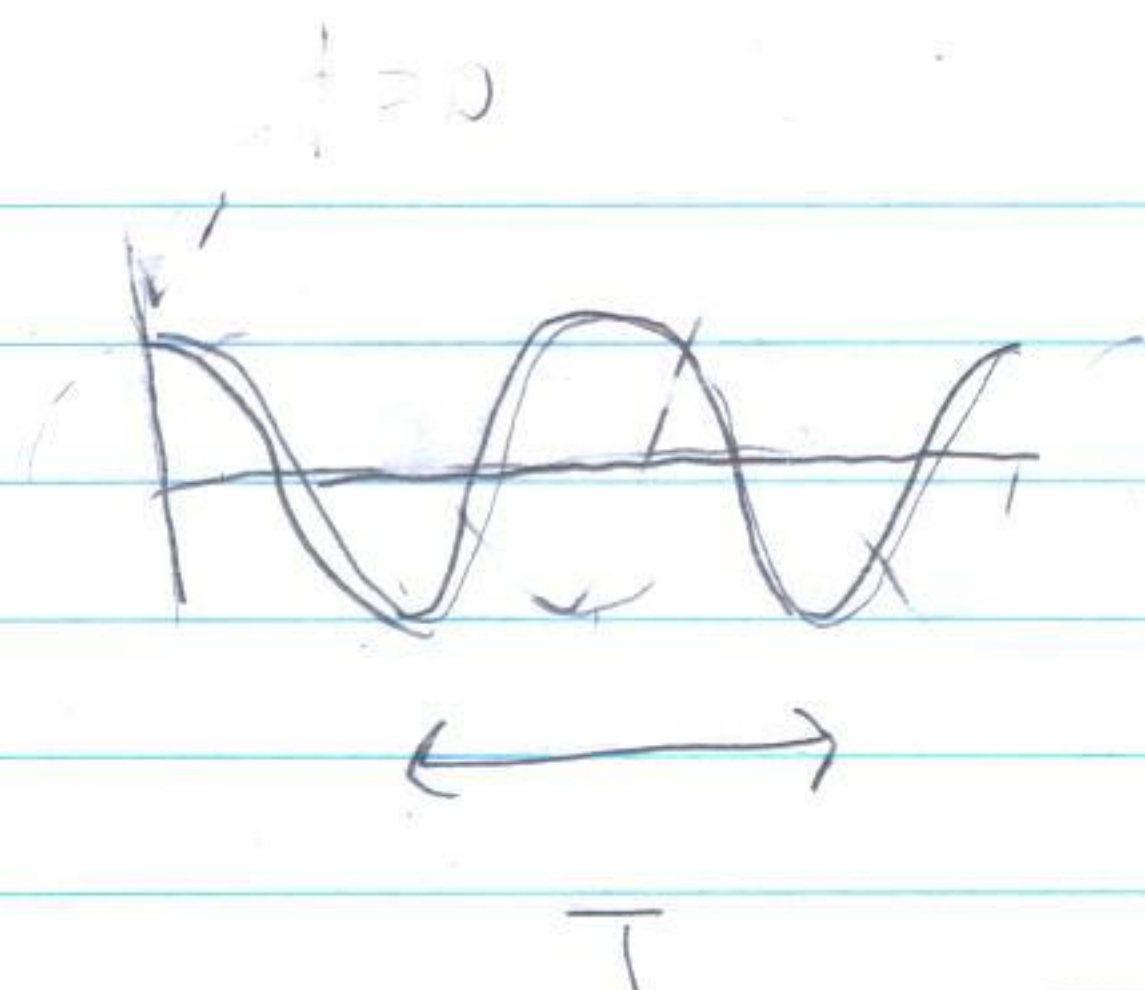
①

What are the constants

$$x = A \cos(\omega_0 t + \phi)$$



②



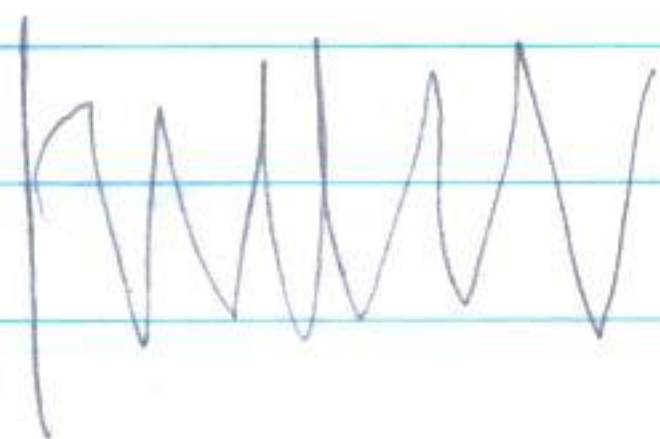
$T = \text{time for one cycle}$

$$(\omega(t+T) + \phi) - (\omega(t) + \phi) = 2\pi$$

$$\omega T = 2\pi$$

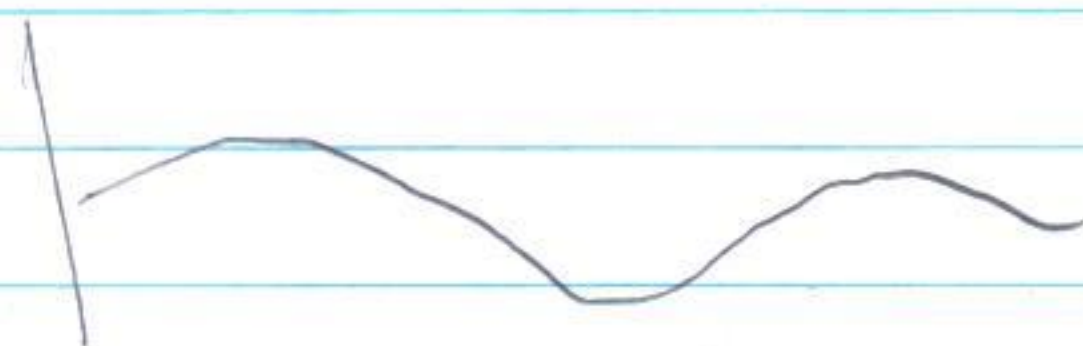
$$T = \frac{2\pi}{\omega} \rightarrow f = \frac{1}{T} = \frac{\text{cycles}}{\text{sec}}$$

Example ① ω large

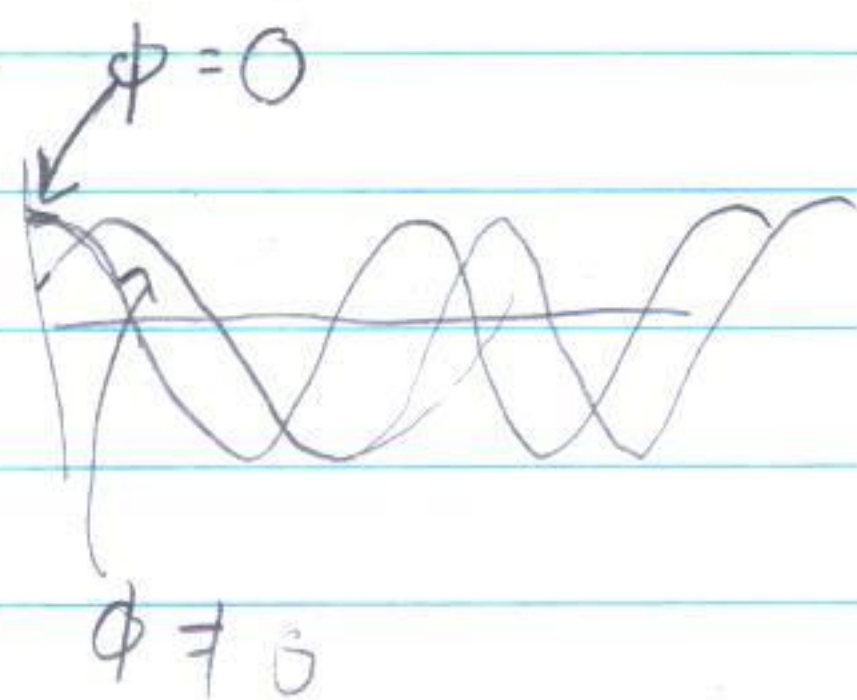


② ω small

= units
Hz

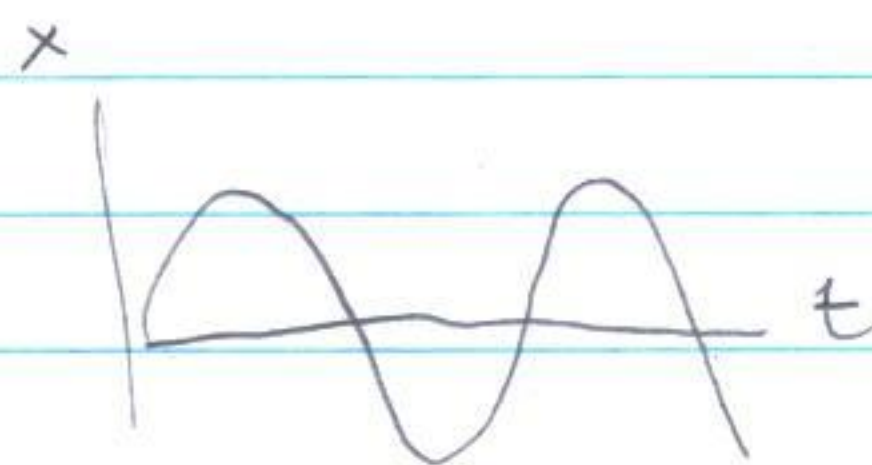


③ What is ϕ ?



$$\phi = \frac{\pi}{2} \quad x = A \cos(\omega t + \frac{\pi}{2})$$

$$x = A \sin(\omega t)$$



Example

A 200g block is connected to a light spring force constant 5.00 N/m and oscillates on a frictionless surface. It is displaced 5.0cm from equilibrium.

a) Determine its period -- time to return

$$\omega = \sqrt{\frac{k}{m}} = \sqrt{\frac{5.0 \text{ N/m}}{0.2 \text{ kg}}} = 5.0 \text{ 1/s}$$

$$T = \frac{2\pi}{\omega} = \frac{2\pi}{5.0 \text{ 1/s}} = 1.26 \text{ s}$$

b) Determine the Amplitude and phase

$$x = A \cos(\omega t + \phi)$$

$$v = \frac{dx}{dt} = -\omega A \sin(\omega t + \phi)$$

we adjust A and ϕ to match the state at $t=0$, called initial condition.

$$x(t=0) = 5 \text{ cm} = A \cos(\omega \cdot 0 + \phi) =$$

$$v(t=0) = 0 = -\omega A \sin(\phi)$$

$$\phi = 0 \quad A = 5 \text{ cm}$$

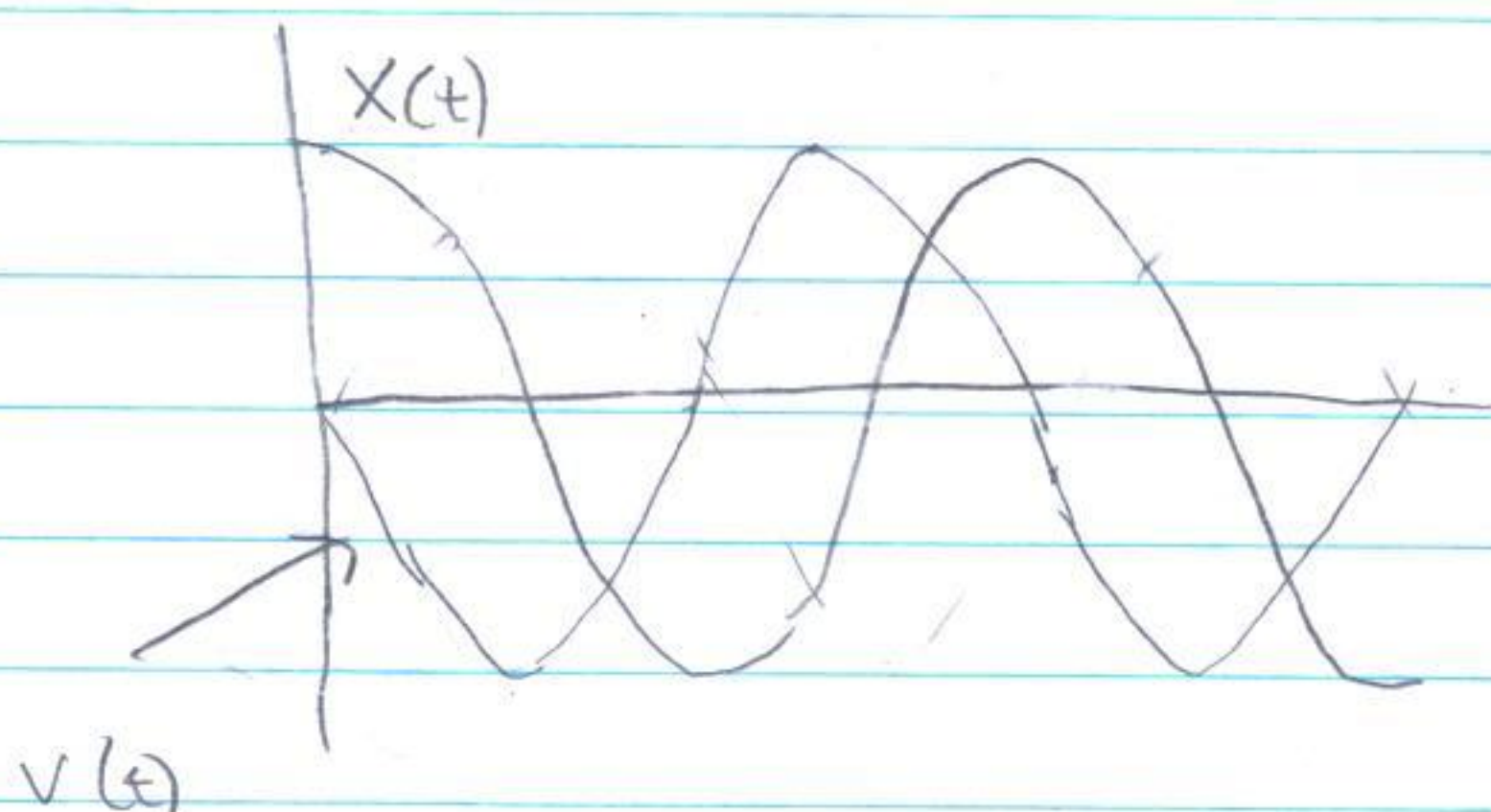
c) Find the position and velocity as a function of time

$$x = A \cos(\omega t + \phi)$$

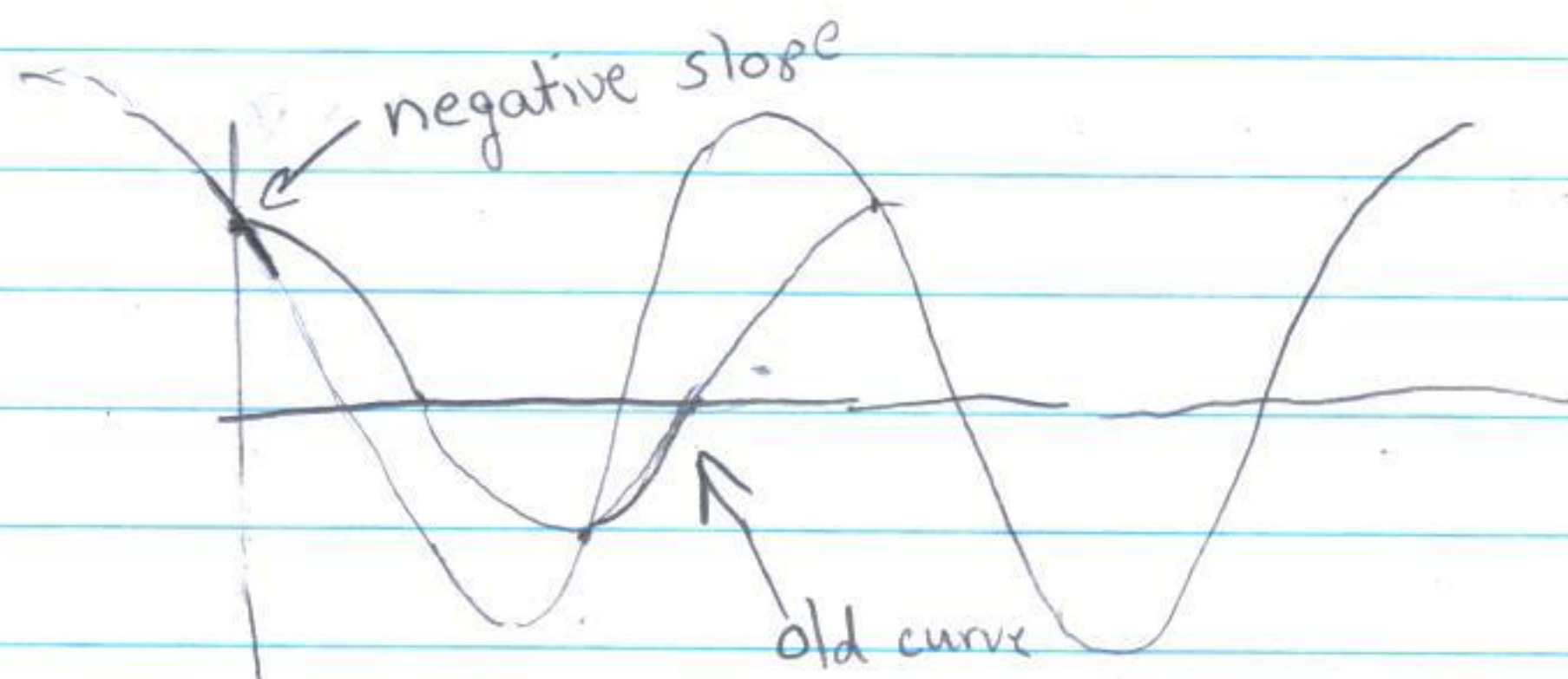
$$x(t) = (0.05 \text{ m}) \cos\left(\frac{5.0}{\text{s}} t\right)$$

$$x_{\text{max}} = 0.05 \text{ m}$$

$$v(t) = -0.25 \frac{\text{m}}{\text{s}} \sin\left(\frac{5.0}{\text{s}} t\right) \rightarrow v_{\text{max}} = -0.25 \text{ m/s}$$



d) Suppose that we give the same block $x_i = 5 \text{ cm}$ a kick back towards equilibrium, $v_i = -0.1 \text{ m/s}$



ϕ and A change but $\omega_0 = \sqrt{\frac{k}{m}}$ does not!

$$x(t) = A \cos(\omega t + \phi)$$

$$v(t) = -A\omega \sin(\omega t + \phi)$$

$$x(0) = A \cos \phi = x_i = 5 \text{ cm}$$

$$v(0) = -A\omega \sin \phi = v_i = -0.1 \text{ m/s}$$

$$\frac{-\omega A \sin \phi}{A \cos \phi} = \frac{v_i}{x_i}$$

$$\tan \phi = \frac{-v_i}{\omega x_i} = - \frac{(-0.1 \text{ m})}{(5.0 \frac{\text{m}}{\text{s}})(0.05 \text{ m})} = 0.4$$

$$\phi = 0.377 \text{ rad} = 0.$$

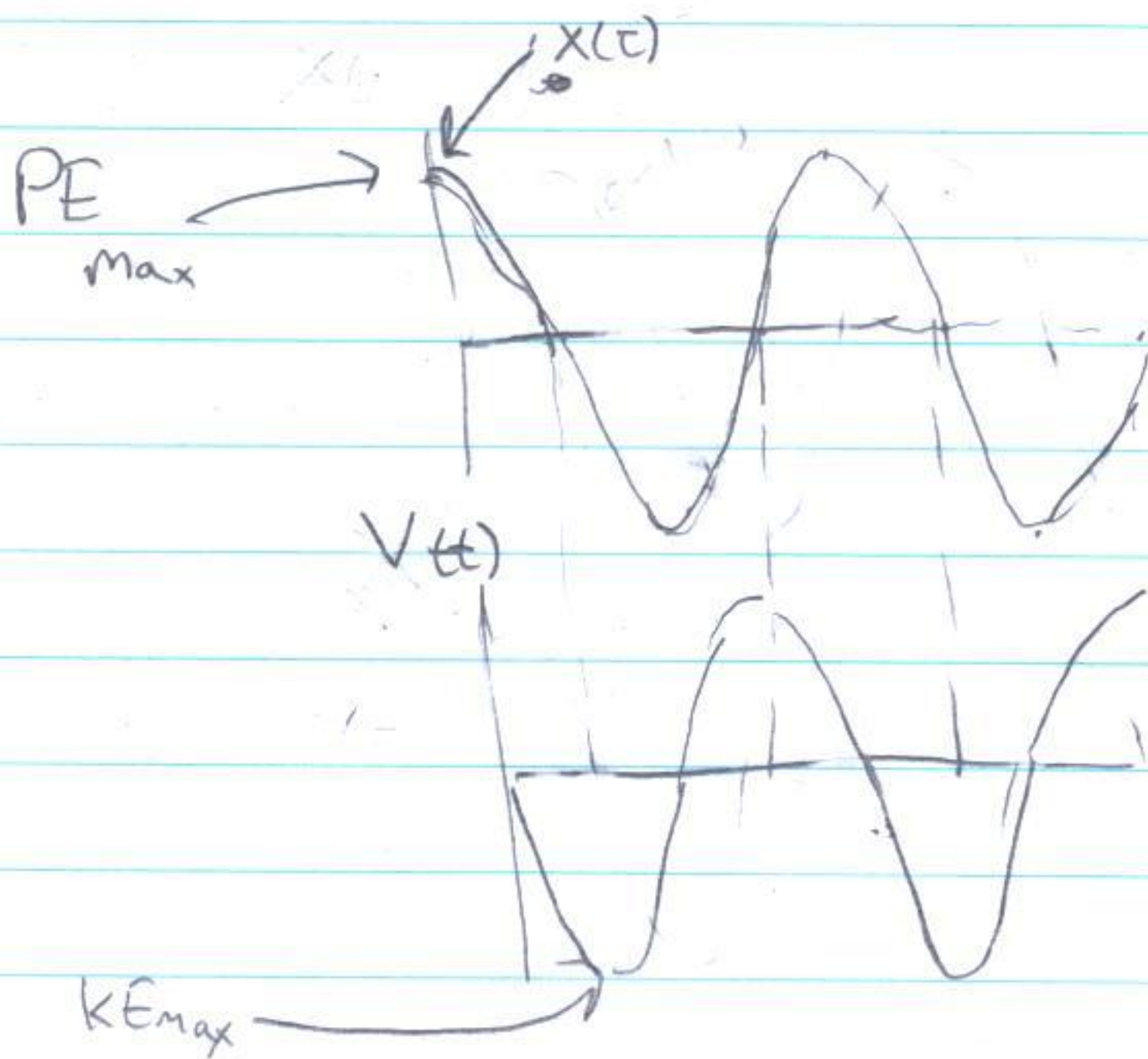
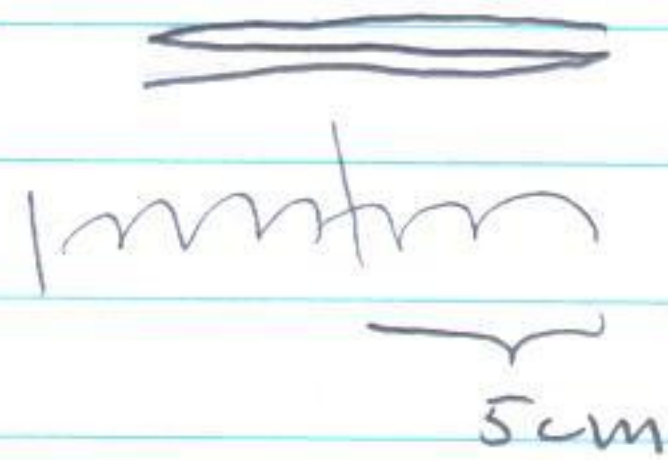
$$A = \frac{x_i}{\cos \phi} = 0.054 \text{ m}$$

A is larger,

Intuitively this is because we have given more energy to the system by kicking it

Energy is SHM.

$$K = \frac{1}{2} m v^2(t)$$



$$x = A \cos \omega t$$

$$v = -A\omega \sin \omega t$$

$$\text{Kinetic} + \text{Potential} = \text{Constant} = \frac{1}{2} k A^2$$

Proof: $x = A \cos \omega t$

$$v = -A\omega \sin \omega t$$

↑
potential
energy
at max
extension

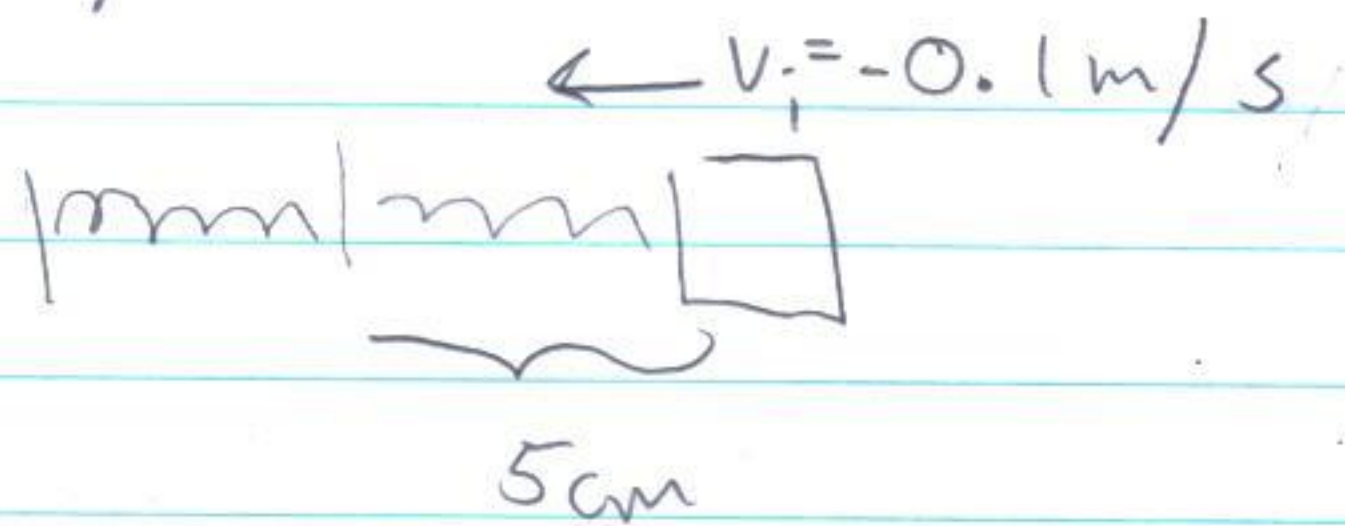
$$\frac{1}{2} k x^2 + \frac{1}{2} m v^2 =$$

$$\frac{1}{2} k A^2 \cos^2 \omega t + \frac{1}{2} m (-A\omega)^2 \sin^2 \omega t$$

$$\frac{1}{2} k A^2 \cos^2 \omega t + \frac{1}{2} \underbrace{k A^2}_{m\omega^2} \sin^2 \omega t$$

$$= \frac{1}{2} k A^2 [\cos^2 \omega t + \sin^2 \omega t] = \frac{1}{2} k A^2$$

Problem, same as before:



Determine the amplitude of oscillation:

Sol

$$\frac{1}{2} k x_i^2 + \frac{1}{2} m v_i^2 = \frac{1}{2} k A^2$$

Energy
at starting
point

Energy at maximum stretch
is all potential $A = x_{\text{max}}$

$$\sqrt{\left(x_i^2 + \frac{m v_i^2}{k} \right)} = A$$

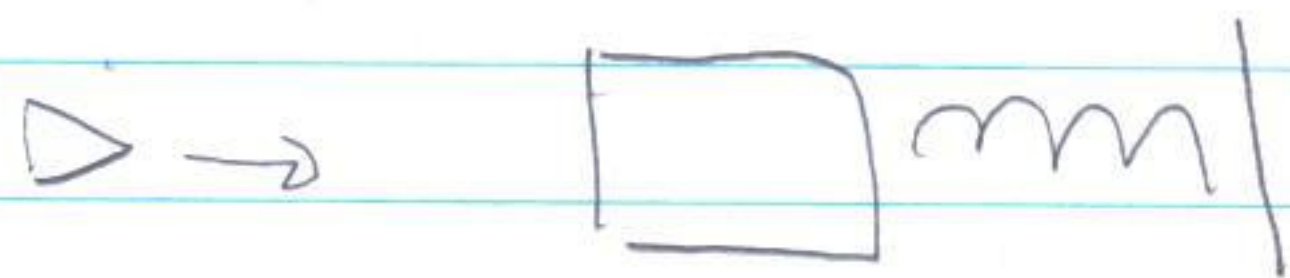
$$\sqrt{\left(x_i^2 + \frac{v_i^2}{\omega_0^2} \right)} = A$$

$$\sqrt{\left((0.05 \text{ m})^2 + \left(\frac{-0.1 \text{ m/s}}{5 \text{ 1/s}} \right)^2 \right)} = A$$

$$\boxed{0.54 \text{ m} = A} \quad \leftarrow \text{same}$$

Make up numbers please
Example ← mass m_B , $v = v_B$ I was in a rush!

A bullet \uparrow is fired into a block mass M attached to a spring ω spring constant k



a) Find the max, compression

b) Find the equation of motion, determine the time it takes to reach $\frac{1}{4}$ of its max amplitude

a) Mom $m_B v_B = (M + m_B) v_0$

$$\frac{m_B v_B}{M + m_B} = v_0$$

$$\frac{1}{2} k A^2 = \frac{1}{2} (M + m_B) v_0^2 = \frac{1}{2} m_B v_B^2 \left(\frac{m_B}{M + m_B} \right)$$

$$A = \left(\frac{M + m_B}{k} \right)^{1/2} v_0^2 = \left(\frac{v_0^2}{\omega_0^2} \right)^{1/2}$$

b) $x(t) = A \sin(\omega t) =$ use rad!

$$\frac{1}{4} A = A \sin(\omega t) \rightarrow t = \sin^{-1}\left(\frac{1}{4}\right) \cdot \sqrt{\frac{M + m_B}{k}}$$

Determine the maximum velocity:

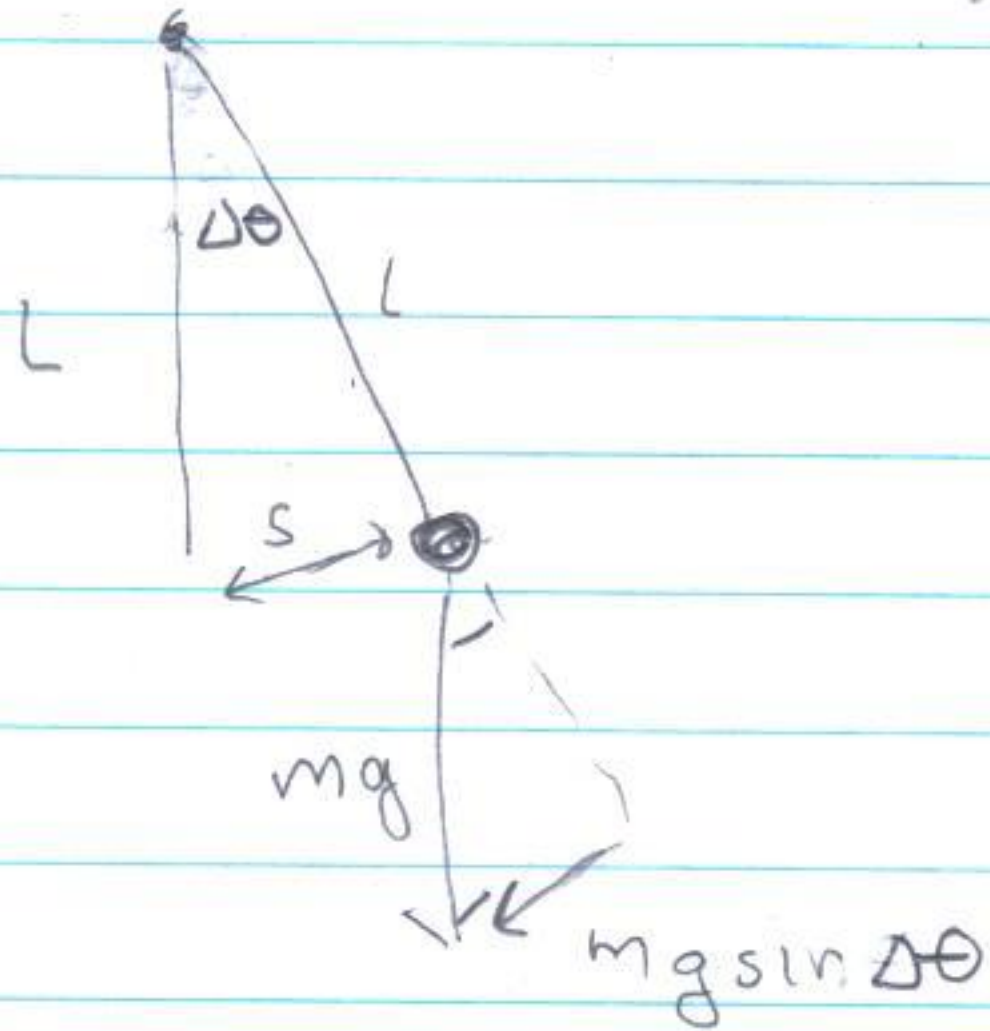
$$\frac{1}{2} k x_i^2 + \frac{1}{2} m v_i^2 = \frac{1}{2} m v_{\max}^2$$

Then there is
no PE

$$\sqrt{\frac{k}{m} x_i^2 + v_i^2} = v_{\max}$$

$$x = A \cos(\omega t + \phi)$$

Pendulum Pendulum



$$s \approx L \Delta\theta$$

$$F = ma$$

$$-mg \sin \Delta\theta = m \frac{d^2 s}{dt^2} = m L \frac{d^2 (\Delta\theta)}{dt^2}$$

$$\sin \Delta\theta \approx \Delta\theta$$

$$-\frac{g}{L} \Delta\theta = \frac{d^2 (\Delta\theta)}{dt^2}$$

$$-\left(\frac{g}{L}\right) \Delta\theta = \frac{d^2 (\Delta\theta)}{dt^2}$$

$$\frac{d^2 x}{dt^2} = -\left(\frac{k}{m}\right) x$$

$$\Theta = A \cos(\omega_0 t + \phi)$$

$$T = \frac{2\pi}{\omega} \text{ etc}$$

$$\omega_0 = \sqrt{\frac{g}{L}}$$


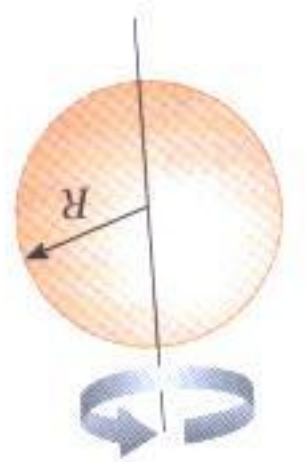


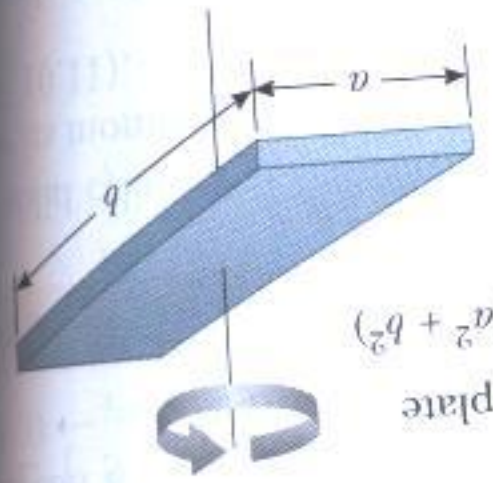
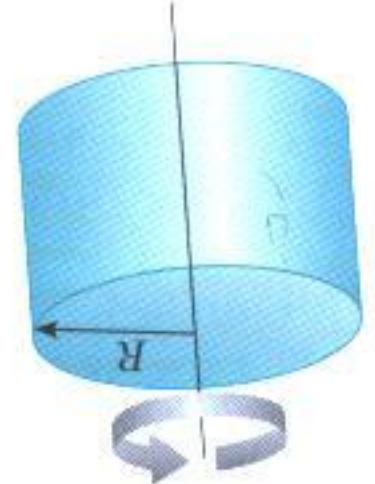
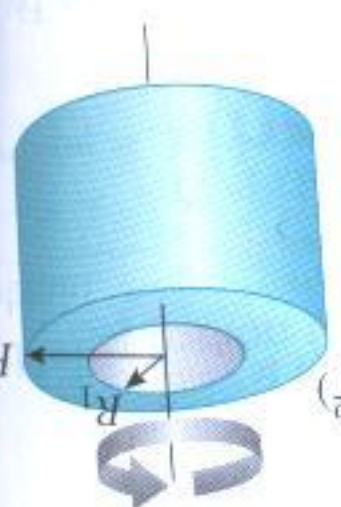
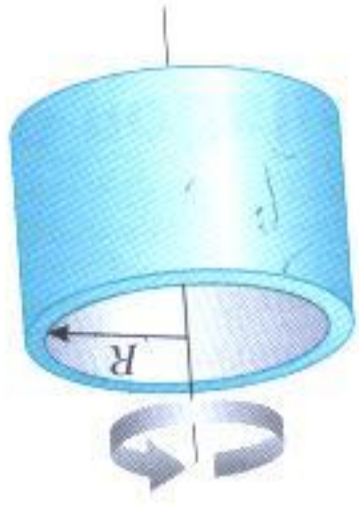
Answer Note that the result for the moment of inertia of a cylinder does not depend on L , the length of the cylinder. In other words, it applies equally well to a long cylinder and a flat disk having the same mass M and radius R . Thus, the moment of inertia of the cylinder would not be affected by changing its length.

Parallel-axis theorem

Table 10.2 gives the moments of inertia for a number of objects about specific axes. The moments of inertia of rigid objects with simple geometry (high symmetry) are relatively easy to calculate provided the rotation axis coincides with an axis of symmetry. However, even for a highly symmetric object, use of an important theorem called the **parallel-axis theorem**, often simplifies the calculation. Suppose the moment of inertia about an axis through the center of mass of an object is I_{CM} and the moment of inertia about an axis parallel to it and a distance D away from this axis is I . To prove the parallel-axis theorem, suppose that an object rotates in the xy -plane about the z axis, as shown in Figure 10.12, and that the coordinates of the center of mass are x_{CM} , y_{CM} . Let the mass element dm have coordinates x , y . Because

$$I = I_{CM} + MD^2$$

Table 10.2 Moments of Inertia of Homogeneous Rigid Objects with Different Geometries

	Thin spherical shell $I_{CM} = \frac{2}{3} MR^2$		Solid sphere $I_{CM} = \frac{2}{5} MR^2$
	Long thin rod with rotation axis through end $I = \frac{1}{3} ML^2$		Long thin rod with rotation axis through center $I_{CM} = \frac{1}{12} ML^2$
	Rectangular plate $I_{CM} = \frac{1}{12} M(a^2 + b^2)$		Solid cylinder or disk $I_{CM} = \frac{1}{2} MR^2$
	Hollow cylinder $I_{CM} = \frac{1}{2} M(R_1^2 + R_2^2)$		Hoop or thin cylindrical shell $I_{CM} = MR^2$